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Search for electroweak production of charginos and neutralinos in WH events in proton-proton collisions at $\sqrt{s} = 13$ TeV

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Search for electroweak production of charginos and neutralinos in WH events in proton-proton collisions at $\sqrt{s} = 13$ TeV



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KEYWORDS: Hadron-Hadron scattering (experiments), Supersymmetry

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1 Introduction

Supersymmetry (SUSY) [1–8] is a theoretically attractive extension of the standard model (SM) that is based on a symmetry between bosons and fermions. SUSY predicts the existence of a superpartner for every SM particle, with the same gauge quantum numbers but differing by one half unit of spin. In R-parity conserving SUSY models, supersymmetric particles are created in pairs, and the lightest supersymmetric particle (LSP) is stable [9–11]. As a result, SUSY also provides a potential connection to cosmology as the LSP, if neutral and stable, may be a viable dark matter candidate.

Previous searches based on 13 TeV proton-proton collision data at the CERN LHC focused on strong production of colored SUSY particles [12–28]. Pair production of these particles would have the largest cross section for SUSY processes and therefore provides the strongest discovery potential with small datasets. However, the absence of signals in these searches suggests that strongly produced SUSY particles may be too massive to be found with the present data. In contrast, neutralinos ($\tilde{\chi}^0$) and charginos ($\tilde{\chi}^\pm$), mixtures of the superpartners of the SM electroweak gauge bosons and the Higgs bosons, can have masses within the accessible range. Because of the absence of color charge, the production cross

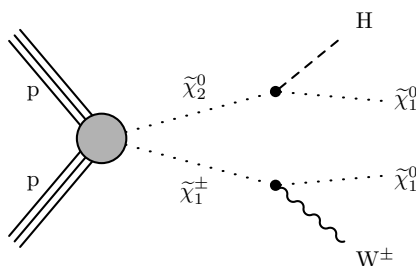


Figure 1. Diagram corresponding to the SUSY simplified model targeted by this analysis, i.e., chargino-neutralino production, with the chargino decaying to a W boson and an LSP, while the heavier neutralino decays to a Higgs boson and an LSP.

sections are lower, and these particles may have thus far eluded detection. This provides strong motivation for dedicated searches for electroweak SUSY particle production.

Depending on the mass spectrum, the charginos and neutralinos can have significant decay branching fractions to vector bosons V (W or Z) and the Higgs boson (H). Here, “ H ” refers to the 125 GeV Higgs boson [29], interpreted as the lightest CP-even state of an extended Higgs sector. The H boson is expected to have SM-like properties if all of the other Higgs bosons are much heavier [30]. The observation of a Higgs boson in a SUSY-like process would provide evidence that SUSY particles couple to the Higgs field, a necessary condition for SUSY to stabilize the Higgs boson mass. Pair production of neutralinos and/or charginos can thus lead to the HH , VH , and VV decay modes, with a large fraction of the possible final states containing at least one isolated lepton. Such events can be easily selected with simple triggers and do not suffer from large quantum chromodynamics multijet background.

In this paper we focus on a simplified model [31–35] of supersymmetric chargino-neutralino ($\tilde{\chi}_1^\pm \tilde{\chi}_2^0$) production with the decays $\tilde{\chi}_1^\pm \rightarrow W^\pm \tilde{\chi}_1^0$ and $\tilde{\chi}_2^0 \rightarrow H \tilde{\chi}_1^0$, as shown in figure 1. Both the $\tilde{\chi}_1^\pm$ and $\tilde{\chi}_2^0$ are assumed to be wino-like and have the same mass. The lightest neutralino $\tilde{\chi}_1^0$, produced in the decays of $\tilde{\chi}_1^\pm$ or the $\tilde{\chi}_2^0$, is considered to be the stable LSP, which escapes detection. When the W boson decays leptonically, this process typically results in a signature with one lepton, two jets that originate from the decay $H \rightarrow b\bar{b}$, and large missing transverse momentum from the neutrino in the W boson decay and the LSPs.

Results of searches for electroweak pair production of SUSY particles were previously reported by the ATLAS and CMS Collaborations using data sets of 8 TeV proton-proton (pp) collisions [36–38] in a variety of event topologies and final states. No excesses above the SM expectations were observed, and the results of those searches were used to place lower limits on the mass of pair-produced charginos and neutralinos. Assuming mass-degenerate $\tilde{\chi}_1^\pm$ and $\tilde{\chi}_2^0$, and sleptons (the SUSY partners of the SM leptons) with lower masses, the searches probed masses up to approximately 700 GeV. For the WH decays assumed here, the strongest mass limit was around 270 GeV. With the increase of the LHC collision energy from 8 to 13 TeV, and a significantly larger data set, searches based on 13 TeV data have the potential to quickly surpass the sensitivity of the previous analyses.

This paper presents the result of a search using a data set corresponding to an integrated luminosity of 35.9 fb^{-1} of pp collisions collected at a center-of-mass energy of 13 TeV with the CMS detector in 2016. The results are interpreted in the simplified SUSY model with chargino-neutralino production depicted in figure 1.

2 The CMS detector

The central feature of the CMS apparatus is a superconducting solenoid, 13 m in length and 6 m in diameter, which provides an axial magnetic field of 3.8 T. Within the field volume are several particle detection systems. Charged-particle trajectories are measured with silicon pixel and strip trackers, covering $0 \leq \phi < 2\pi$ in azimuth and $|\eta| < 2.5$ in pseudorapidity, where $\eta \equiv -\ln[\tan(\theta/2)]$ and θ is the polar angle of the trajectory of the particle with respect to the counterclockwise beam direction. The transverse momentum, the component of the momentum p in the plane orthogonal to the beam, is defined in terms of the polar angle as $p_T = p \sin \theta$. A lead-tungstate crystal electromagnetic calorimeter and a brass and scintillator hadron calorimeter surround the tracking volume, providing energy measurements of electrons, photons, and hadronic jets in the range $|\eta| < 3.0$. Muons are identified and measured within $|\eta| < 2.4$ by gas-ionization detectors embedded in the steel flux-return yoke of the solenoid. Forward calorimeters on each side of the interaction point encompass $3.0 < |\eta| < 5.0$. The detector is nearly hermetic, allowing momentum imbalance measurements in the plane transverse to the beam direction. A two-tier trigger system selects pp collision events of interest for use in physics analyses. A detailed description of the CMS detector, together with a definition of the coordinate system used and the relevant kinematic variables, can be found in ref. [39].

3 Event samples, reconstruction, and selection

3.1 Object definition and preselection

Event reconstruction is based on the particle-flow (PF) algorithm [40, 41], which combines information from the tracker, calorimeter, and muon systems to reconstruct and identify PF candidates, i.e., charged and neutral hadrons, photons, muons, and electrons. To select collision events, we require at least one reconstructed vertex. The reconstructed vertex with the largest value of summed physics-object p_T^2 is taken to be the primary pp interaction vertex. The physics objects are the objects returned by a jet finding algorithm [42, 43] applied to all charged tracks associated with the vertex, plus the corresponding associated missing transverse momentum. The missing transverse momentum vector, \vec{p}_T^{miss} , is defined as the negative vector sum of the momenta of all reconstructed PF candidates projected onto the plane perpendicular to the proton beams. Its magnitude is referred to as E_T^{miss} . Events with possible contributions from beam halo processes or anomalous noise in the calorimeter can have large values of E_T^{miss} and are rejected using dedicated filters [44].

Data events are selected using triggers that require the presence of an isolated electron or muon with p_T thresholds of 27 GeV or 24 GeV, respectively. Muon events may also be accepted using a trigger that does not require isolation but instead requires $p_T > 50$ GeV.

The trigger efficiency, measured using a data sample of $Z/\gamma^* \rightarrow \ell\ell$ events, varies in the range 70–95% (85–92%) depending on the η and p_T of the electron (muon).

Selected events are required to have exactly one lepton (electron or muon), with electrons (muons) satisfying $p_T > 30(25)$ GeV and $|\eta| < 1.44(2.1)$. Electron candidates are reconstructed starting from a cluster of energy deposits in the electromagnetic calorimeter. The cluster is then matched to a reconstructed track. The electron selection is based on the shower shape, track-cluster matching, and consistency between the cluster energy and the track momentum [45]. Muon candidates are reconstructed by performing a global fit that requires consistent hit patterns in the tracker and the muon system [46]. For both lepton flavors, the impact parameter with respect to the primary vertex is required to be less than 0.5 mm in the transverse plane and 1 mm along the beam direction.

Leptons are required to be isolated from other activity in the event. A measure of lepton isolation is the scalar p_T sum (p_T^{sum}) of all PF candidates not associated with the lepton within a cone of radius $\Delta R \equiv \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2} = 0.3$, where $\Delta\eta$ and $\Delta\phi$ are the distances between the lepton and the PF candidates at the primary vertex in η - ϕ space [47]. Only charged PF candidates compatible with the primary vertex are included in the sum. The average contribution of particles from additional pp interactions in the same or nearby bunch crossings (pileup) is subtracted from p_T^{sum} . We require that p_T^{sum} be less than 5 GeV. Typical lepton identification and isolation efficiencies, measured in samples of $Z/\gamma^* \rightarrow \ell\ell$ events, are approximately 80–85% (85–90%) for electrons (muons), depending on p_T and η .

Particle-flow candidates are clustered to form jets using the anti- k_T clustering algorithm [42] with a distance parameter of 0.4, as implemented in the FASTJET package [43]. Only charged PF candidates compatible with the primary vertex are used in the clustering. The pileup contribution to the jet energy is estimated on an event-by-event basis using the jet area method described in [48] and is subtracted from the overall jet p_T . Corrections are applied to the energy measurements of jets to account for non-uniform detector response and are propagated consistently as a correction to \vec{p}_T^{miss} [49, 50]. The selected lepton can also be reconstructed as a jet, so any jets within $\Delta R = 0.4$ of the lepton are removed from the list of considered jets.

Selected events are required to contain exactly two jets with $p_T > 30$ GeV and $|\eta| < 2.4$. Both of these jets must be consistent with containing the decay products of a heavy-flavor (HF) hadron, as identified using the combined secondary vertex (CSVv2) tagging algorithm [51]. Such jets are referred to as b jets. The CSVv2 algorithm has three main operating points: loose, medium, and tight. We require both jets to be tagged according to the loose operating point, and at least one of them to be tagged with the medium operating point. The efficiency of this algorithm for jets arising from b quarks with p_T between 30 and 400 GeV is in the range 60–65% (70–75%) for the medium (loose) working point. The misidentification rate for jets arising from light quarks or gluons is approximately 1% (10%) for the medium (loose) working point.

The largest background in this search arises from $t\bar{t}$ and tW events with decays into two-lepton final states in which one of the leptons is not reconstructed or identified. In order to reduce these backgrounds, we search for a second electron or muon with $p_T > 5$ GeV and looser identification and isolation requirements, and reject events where such a lepton

is found. Second leptons are required to satisfy $p_T^{\text{sum}}/p_T < 0.1$, where p_T^{sum} is calculated here with a cone radius of $\Delta R = 0.2$ for $p_T^{\text{lep}} \leq 50$ GeV, and $\Delta R = \max(0.05, 10 \text{ GeV}/p_T^{\text{lep}})$ at higher values of lepton transverse momentum. We also reject events with reconstructed hadronically decaying tau leptons with $p_T > 20$ GeV [52], or isolated tracks with $p_T > 10$ GeV and opposite electric charge relative to the selected lepton. For this purpose, a track is considered isolated if $p_T^{\text{sum}}/p_T < 0.1$ and $p_T^{\text{sum}} < 6$ GeV, where p_T^{sum} here is constructed with charged PF candidates compatible with the primary vertex, the cone radius is $\Delta R = 0.3$, and p_T is the transverse momentum of the track.

The final two requirements that complete the preselection are $E_T^{\text{miss}} \geq 125$ GeV and $M_T > 50$ GeV, where M_T is the transverse mass of the lepton- E_T^{miss} system, defined as

$$M_T = \sqrt{2p_T^{\text{lep}} E_T^{\text{miss}} [1 - \cos(\Delta\phi)]}, \quad (3.1)$$

where p_T^{lep} is the transverse momentum of the lepton and $\Delta\phi$ is the angle between the transverse momentum of the lepton and \vec{p}_T^{miss} .

3.2 Signal region definition

The signal regions are defined by additional requirements on the kinematic properties of preselected events. The invariant mass of the two b jets is required to be in the range $90 \leq M_{b\bar{b}} \leq 150$ GeV, consistent with the Higgs boson mass within the resolution. The $M_{b\bar{b}}$ distribution for signal and background processes is shown in figure 2 (top left), displaying a clear peak for signal events near the Higgs boson mass.

To suppress single-lepton backgrounds originating from semileptonic $t\bar{t}$, W + jets, and single top quark processes, the preselection requirement on M_T is tightened to >150 GeV. This is because the M_T distribution in these processes with a single leptonically decaying W boson has a kinematic endpoint $M_T < m_W$, where m_W is the W boson mass. The endpoint can be exceeded by off mass-shell W bosons or because of detector resolution effects. The M_T requirement significantly reduces single-lepton backgrounds, as shown in figure 2 (bottom left).

In order to further suppress both semileptonic and dileptonic $t\bar{t}$ backgrounds, we utilize the contranverse mass variable, M_{CT} [53, 54]:

$$M_{CT} = \sqrt{2p_T^{b1} p_T^{b2} [1 + \cos(\Delta\phi_{bb})]}, \quad (3.2)$$

where p_T^{b1} and p_T^{b2} are the transverse momenta of the two jets, and $\Delta\phi_{bb}$ is the azimuthal angle between the pair. As shown in refs. [53, 54], this variable has a kinematic endpoint at $(m^2(\delta) - m^2(\alpha))/m(\delta)$, where δ is the pair-produced heavy particle and α is the invisible particle produced in the decay of δ . In the case of $t\bar{t}$ events, when both jets from b quarks are correctly identified, the kinematic endpoint corresponds to the top quark mass, while signal events tend to have higher values of M_{CT} . This is shown in figure 2 (bottom right). We require $M_{CT} > 170$ GeV.

After all other selections, we define two exclusive bins in E_T^{miss} to enhance sensitivity to signal models with different mass spectra: $125 \leq E_T^{\text{miss}} < 200$ GeV and $E_T^{\text{miss}} \geq 200$ GeV. The E_T^{miss} distribution is shown in figure 2 (top right).

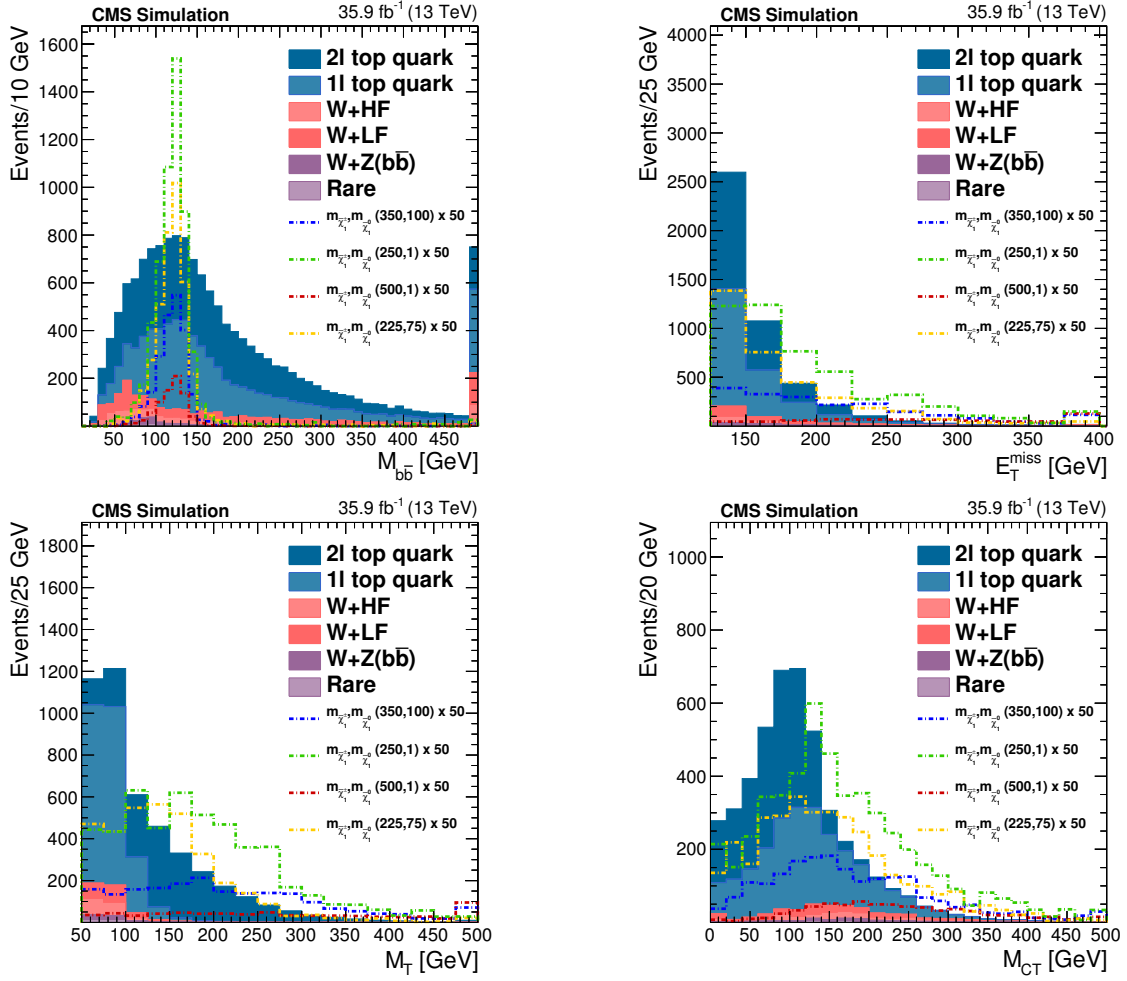


Figure 2. Distributions in $M_{b\bar{b}}$ (top left), E_T^{miss} (top right), M_T (bottom left), and M_{CT} (bottom right) for signal and background events in simulation after the preselection. The E_T^{miss} , M_T , and M_{CT} distributions are shown after the $90 < M_{b\bar{b}} < 150$ GeV requirement. Expected signal distributions are also overlaid as open histograms for various mass points, with the signal cross section scaled up by a factor of 50 for display purposes. The legend entries for signal give the masses $(m_{\tilde{\chi}_1^\pm}, m_{\tilde{\chi}_1^0})$ in GeV and the factor by which the signal cross section has been scaled.

3.3 Signal and background simulation

Samples of $t\bar{t}$, W + jets, and Z + jets events, as well as $t\bar{t}$ production in association with a vector boson, are generated using the MADGRAPH5_aMC@NLO 2.2.2 [55] generator at leading order (LO) with the MLM matching scheme [56], while tW and single top quark t -channel events are generated at next-to-leading-order (NLO) using POWHEG V2 [57–59]. A top quark mass of $m_t = 172.5$ GeV, and the NNPDF3.0 LO or NLO [60] parton distribution functions (PDFs) are used in the event generation. Single top quark s -channel production is simulated using MADGRAPH5_aMC@NLO 2.2.2 at NLO precision with the FxFx matching scheme [61]. Samples of diboson (WW , WZ , and ZZ) events are generated with either POWHEG or MADGRAPH5_aMC@NLO at NLO precision. Normalization of the

simulated background samples is performed using the most accurate cross section calculations available [55, 62–72], which generally correspond to NLO or next-to-NLO precision.

The chargino-neutralino signal samples are also generated with MADGRAPH5_aMC@NLO at LO precision. For these samples we improve on the modeling of initial-state radiation (ISR), which affects the total transverse momentum (p_T^{ISR}) of the system of SUSY particles, by reweighting the p_T^{ISR} distribution in these events. This reweighting procedure is based on studies of the p_T of Z bosons [73]. The reweighting factors range between 1.18 at $p_T^{\text{ISR}} = 125$ GeV and 0.78 for $p_T^{\text{ISR}} > 600$ GeV. We take the deviation from 1.0 as the systematic uncertainty in the reweighting procedure.

Parton showering and fragmentation in all of these samples are performed using PYTHIA V8.212 [74] with the CUETP8M1 tune [75]. For both signal and background events, additional simultaneous proton-proton interactions (pileup) are generated with PYTHIA and superimposed on the hard collisions. The response of the CMS detector for SM background samples is simulated using GEANT4-based model [76], while that for new physics signals is performed using the CMS fast simulation package [77]. All simulated events are processed with the same chain of reconstruction programs as that used for collision data.

Small differences between the b tagging efficiencies measured in data and simulation are corrected using data-to-simulation scale factors. Corrections are also applied to account for differences between lepton selection efficiencies (trigger, reconstruction, identification, and isolation) in data and simulation.

4 Backgrounds

The backgrounds for this search are classified into six categories. The first and most important category, referred to as dilepton top quark events, consists mainly of events from top quark pair production with both quarks decaying leptonically, but also including contributions from the associated production of a single top quark with a W boson, both of which decay leptonically. The second to fifth categories include processes with a single leptonically decaying W boson. Events with a single W are divided into two categories: those with b jets (W+HF, for “heavy flavor”) and those without (W+LF, for “light flavor”). A separate category comprises WZ events in which the Z boson decays to $b\bar{b}$ ($WZ \rightarrow \ell\nu b\bar{b}$). Events with one leptonically decaying top quark, either from $t\bar{t}$ or from single top quark t - or s -channel production, are included in the fifth category (“single-lepton top quark”). Finally, other SM processes contribute a small amount to the expected yield in the signal region and are grouped together in the “rare” category. This includes events from Z + jets, WW, WZ (except the decays described above), ZZ, triboson, $t\bar{t}W$, $t\bar{t}Z$, and $WH \rightarrow \ell\nu b\bar{b}$ processes.

All background processes are modeled using MC simulation. Three data control regions (CRs) are defined by inverting the signal region selection requirements, as summarized in table 1. The CRs are defined at both preselection and signal region selection levels. The CRs at the preselection level are defined with looser cuts in order to check the modeling of key discriminant variables. The CRs after the signal region level selections are used to validate the modeling of the main backgrounds and to assign systematic uncertainties in the background predictions. The regions CRMb \bar{b} and CR0b are split into two bins in E_T^{miss}

Selection	Signal regions	CR2 ℓ	CR0b	CRMb \bar{b}
N(leptons)	=1	=1 or 2	=1	=1
Isolated track veto	✓	inverted if 1 ℓ	✓	✓
Tau candidate veto	✓	inverted if 1 ℓ	✓	✓
Number of b tags	=2	=2	=0	=2
Preselection level				
$M_{b\bar{b}}$ [GeV]	—	—	$\in[90,150]$	$\notin[90,150]$
E_T^{miss} [GeV]	≥ 125	≥ 125	≥ 125	≥ 125
M_T [GeV]	> 50	> 150	> 50	> 150
M_{CT} [GeV]	—	—	> 170	—
Signal region selection level				
$M_{b\bar{b}}$ [GeV]	$\in[90,150]$		$\in[90,150]$	$\notin[90,150]$
E_T^{miss} [GeV]	$[125, 200), \geq 200$	not used	$[125, 200), \geq 200$	$[125, 200), \geq 200$
M_T [GeV]	> 150		> 150	> 150
M_{CT} [GeV]	> 170		> 170	> 170

Table 1. Event selections in signal and control regions. The region CR2 ℓ is only used at the preselection level.

to mirror the signal region selection. The expected signal contribution in any of the CRs is always less than 1% of the total SM yields, and typically much smaller.

The dilepton top quark background can be isolated in the CR2 ℓ control region by selecting dilepton events. In addition to a lepton passing the analysis selections, events must contain one of the following: a second electron or muon, an isolated track candidate, or a tau lepton candidate. The latter categories are included to accept hadronically decaying tau leptons. If all the kinematic selections used for the signal regions are applied, the number of events in CR2 ℓ is too low to validate the modeling of the dilepton top quark background. Therefore, this CR is used primarily to validate the modeling of $M_{b\bar{b}}$.

Since the signal produces a resonant peak in the $M_{b\bar{b}}$ distribution, the requirement on $M_{b\bar{b}}$ is inverted to define the background-dominated control region CRMb \bar{b} , which includes a mixture of all backgrounds in proportions similar to those in the signal region. Consequently, this control region is dominated by the dilepton top quark background and is used to validate the modeling of these processes in the kinematic tails of the E_T^{miss} , M_T , and M_{CT} distributions.

The CR0b region is designed to study the W + LF background. It is used to validate the modeling of the kinematic tails in E_T^{miss} , M_T , and M_{CT} for W + jets processes. In this region, the dijet mass M_{jj} computed from the two selected jets is used in place of $M_{b\bar{b}}$.

The background estimation and the associated uncertainties are described in the following sections.

4.1 Dilepton top quark backgrounds

The dilepton top quark process contributes to the event sample in the signal region when the second lepton is not reconstructed or identified. Due to the presence of two neutrinos,

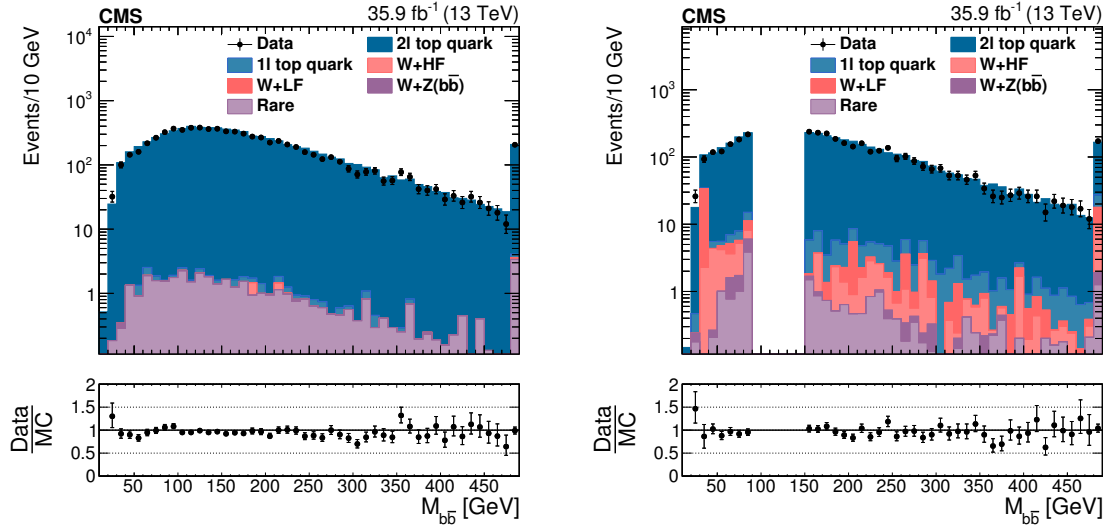


Figure 3. (left) Distribution in $M_{b\bar{b}}$ in CR2 ℓ after the preselection level cuts defined in table 1, comparing data to MC simulation. (right) Distribution in $M_{b\bar{b}}$ in CRMb \bar{b} after preselection level cuts defined in table 1. The signal region range of $90 \leq M_{b\bar{b}} \leq 150$ GeV has been removed from the plot.

these events tend to have higher E_T^{miss} than the single-lepton backgrounds, and their M_T distribution is not bounded by the W boson mass. However, as mentioned above, the M_{CT} requirement significantly suppresses dilepton top quark events. The modeling of this background is validated in two steps. First, the modeling of the $M_{b\bar{b}}$ distribution is validated in CR2 ℓ ; second, the modeling in the kinematic tails of the E_T^{miss} , M_T , and M_{CT} distributions is validated in CRMb \bar{b} . Distributions of $M_{b\bar{b}}$ in CR2 ℓ and CRMb \bar{b} , after the preselection level cuts defined in table 1, are displayed in figure 3 (left) and figure 3 (right), respectively.

In CR2 ℓ , we observe agreement between data and MC, validating the modeling of the $M_{b\bar{b}}$ distribution. We then use CRMb \bar{b} at the signal region selection level to derive a scale factor for the dilepton top quark background separately in each of the analysis E_T^{miss} bins. All other background components are subtracted from the observed data yields, and the result is compared to the dilepton top quark MC prediction. Agreement is observed in the higher E_T^{miss} bin within statistical uncertainties. For the lower E_T^{miss} bin, we find fewer events in data than predicted, and we derive a scale factor of 0.72 for the dilepton top quark background in this bin. From the statistical precision of the data, we assign a systematic uncertainty of 30% in the prediction for both bins. This accounts for any effects that could impact the modeling of this background in simulation, including generator assumptions on factorization and renormalization scales, and PDFs, as well as experimental uncertainties in the jet energy scale, the lepton identification and isolation, trigger, and b tagging efficiencies.

4.2 W boson backgrounds

The M_T requirement (>150 GeV) effectively suppresses the contribution from W + jets events. However, as discussed above, events from W + jets can still enter the M_T tail due

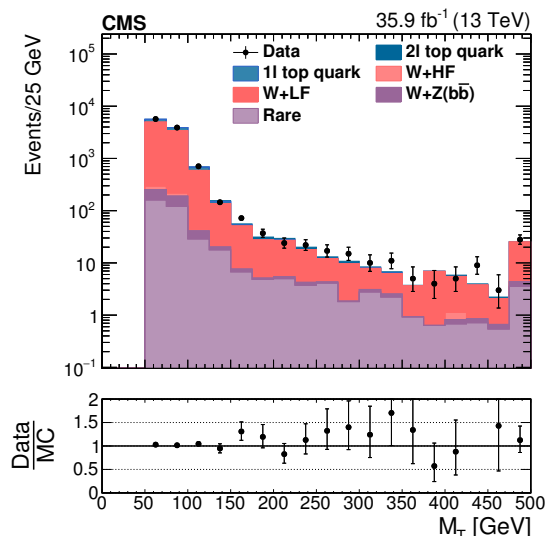


Figure 4. Distribution in M_T in CR0b after the preselection level cuts defined in table 1.

to off-shell W production or E_T^{miss} resolution effects. The control region CR0b consists mostly of W + LF events and is therefore used to validate the modeling of W + jets in the tails of all kinematic variables such as M_T .

Figure 4 shows the M_T distributions of data and simulated events in CR0b after the preselection requirements. The data and simulation agree within uncertainties. The observed yields in data are then compared with MC predictions after applying all the kinematic requirements at signal region selection level defined in table 1. We find agreement within statistical uncertainties. Based on the statistical precision of the data, we assign a 10% systematic uncertainty in the W + jets prediction. This procedure directly tests the W + jets background prediction in the kinematic phase space of the signal region, including experimental uncertainties in the jet energy scale, in the efficiencies for trigger, lepton identification and isolation. It also accounts for most generator assumptions. Additional uncertainties for effects not tested by this procedure are discussed below.

For the W + LF background, the uncertainty due to the b tagging requirements is evaluated by varying the b tagging efficiencies within their measured uncertainties. The uncertainty in the yield in the signal regions is 1%.

For the W + HF background the effects contributing to the kinematic tails are similar to those in W + LF. In this case the tail of the M_T distribution receives contributions from off-shell W boson production and E_T^{miss} resolution effects, but also from neutrinos in semileptonic decays within the b jets. Since this last effect is accounted for in the event generation, we do not apply any additional correction or uncertainty for kinematic tail modeling beyond the one derived above in CR0b.

The most uncertain aspect of the prediction for the W + HF background is the estimate of its cross section relative to the W + LF process. We assign a 50% uncertainty to the normalization of this background [78]. This uncertainty is validated by comparing data to simulation in a CR with $60 < M_T < 120$ GeV and with one or two jets, where the dominant

contribution to the event sample is from $W + \text{jets}$. We find that the 50% uncertainty conservatively covers differences between data and simulation as a function of the number of b jets. Finally, the uncertainty in this prediction due to the uncertainty in the b tagging efficiency is also evaluated and found to be 5%.

The effects discussed above also contribute to the tail of the M_T distribution in $WZ \rightarrow \ell\nu b\bar{b}$ events. As a result, the tail modeling systematic uncertainty for this background is taken to be the same as those evaluated in CR0b. An additional uncertainty of 12% is applied to the normalization of the $WZ \rightarrow \ell\nu b\bar{b}$ background, based on the CMS cross section measurement of inclusive WZ production at 13 TeV [79]. A unique aspect of the $WZ \rightarrow \ell\nu b\bar{b}$ background is that $M_{b\bar{b}}$ peaks at the Z boson mass, at the lower edge of the $M_{b\bar{b}}$ selection used in this analysis. Uncertainties in the jet energy scale can therefore strongly impact the prediction of this background. By varying the jet energy scale within its uncertainty, we derive an uncertainty of 27% in the $WZ \rightarrow \ell\nu b\bar{b}$ background prediction. While this uncertainty is large, the absolute magnitude of this background remains very small in the signal region. Finally, the uncertainty in the background prediction for this process due to the uncertainty in the b tagging efficiency is 2%.

4.3 Other backgrounds

The single-lepton top quark backgrounds are highly suppressed by several of the selections applied in this analysis. Since these contain exactly one leptonically decaying W boson, the M_T requirement is an effective discriminant against them. Requiring exactly two jets also suppresses the $t\bar{t} \rightarrow \ell + \text{jets}$ background, which typically has four jets in the final state. As a result, this background comprises a small fraction of the expected SM prediction in the signal region.

Isolating the single-lepton top quark background in a region kinematically similar to the signal region is difficult since dilepton top quark events tend to dominate when requiring large M_T . The main source of uncertainty in the prediction of this background is the modeling of the E_T^{miss} resolution, which was found to be well modeled in the study of CR0b.

Additional studies of E_T^{miss} resolution are performed using $\gamma + \text{jets}$ events following the method used in ref. [78]. The resolution in data is found to be up to 20% worse than in simulation, leading to higher single-lepton top quark yields than expected from simulation. However, the impact of this effect on the total background prediction is negligible. Due to the difficulties in defining a dedicated control region for this process, we assign a conservative uncertainty of 100% to the single lepton top quark background prediction.

The “rare” backgrounds contribute less than 15% of the expected yield in the signal region. We apply an uncertainty of 50% on the event yields from these processes.

5 Results

Figure 5 shows the distributions of $M_{b\bar{b}}$ in data compared with the SM background prediction after all signal region requirements except the $M_{b\bar{b}}$ selection. No significant deviations from the predictions are observed. Table 2 shows the expected SM background yields in the signal region compared to the observation, as well as predicted yields for several signal

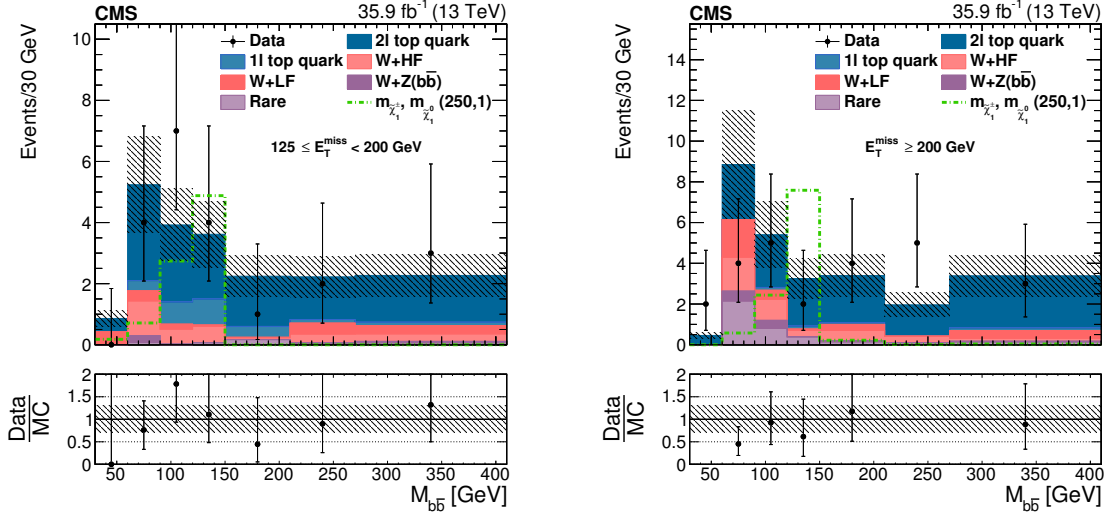


Figure 5. Distributions in $M_{b\bar{b}}$ after all signal region kinematic requirements for the two exclusive E_T^{miss} bins (left: $125 \leq E_T^{\text{miss}} < 200$ GeV, right: $E_T^{\text{miss}} \geq 200$ GeV). The signal region is $90 \leq M_{b\bar{b}} \leq 150$ GeV. The hatched band shows the total uncertainty in the background prediction, including statistical and systematic components. The expected signal distribution for a reference SUSY model is overlaid as an open histogram, and the legend (on the last line) gives the masses as $(m_{\tilde{\chi}_1^\pm}, m_{\tilde{\chi}_1^0})$ in GeV.

models with the masses $(m_{\tilde{\chi}_1^\pm}, m_{\tilde{\chi}_1^0})$ indicated in GeV. The correlation coefficient for the background prediction between the two bins is 0.61. The correlation is incorporated in the likelihood model described below for the interpretation of the results, and it can be used to reinterpret these results in other signal models [80].

6 Interpretation

The results of this analysis are interpreted in the context of the simplified SUSY model depicted in figure 1, $\tilde{\chi}_1^\pm \tilde{\chi}_2^0 \rightarrow W^\pm H \tilde{\chi}_1^0 \tilde{\chi}_1^0$. The $\tilde{\chi}_1^\pm$ and $\tilde{\chi}_2^0$ are assumed to have the same mass, and the branching fractions for the decays listed above are taken to be 100%. The W and Higgs bosons are taken to decay according to their SM branching fractions. Cross section limits as a function of the SUSY particle masses are set using a modified frequentist approach, employing the CL_s criterion and an asymptotic formulation [81–84]. Both signal regions are considered simultaneously in setting limits. The “expected” limit is that under the background-only hypothesis, while the “observed” limit reflects the data yields in the signal regions. The production cross sections are computed at NLO plus next-to-leading-log (NLL) precision in a limit of mass-degenerate wino $\tilde{\chi}_1^\pm$ and $\tilde{\chi}_2^0$, light bino $\tilde{\chi}_1^0$, and with all the other sparticles assumed to be heavy and decoupled [85, 86]. The uncertainty in the cross section calculation includes variations of factorization and renormalization scales, and of the PDFs.

The systematic uncertainties in the signal yield are summarized in table 3. The signal models with the largest acceptance uncertainties are those with $\Delta m = m_{\tilde{\chi}_2^0} - m_{\tilde{\chi}_1^0} \simeq m_H$.

	$125 \leq E_T^{\text{miss}} < 200 \text{ GeV}$	$E_T^{\text{miss}} \geq 200 \text{ GeV}$
Dilepton top quark	4.6 ± 1.5	4.9 ± 1.7
W + LF	0.2 ± 0.1	0.5 ± 0.4
W + HF	1.0 ± 0.9	1.3 ± 1.0
WZ $\rightarrow \ell\nu b\bar{b}$	0.1 ± 0.1	0.4 ± 0.2
Single-lepton top quark	1.6 ± 1.6	0.3 ± 0.4
Rare	$0.0^{+0.2}_{-0.0}$	1.2 ± 0.7
Total SM background	7.5 ± 2.5	8.7 ± 2.2
Data	11	7
$\tilde{\chi}_1^\pm \tilde{\chi}_1^0$ (225,75)	2.4 ± 0.4	2.3 ± 0.4
$\tilde{\chi}_1^\pm \tilde{\chi}_1^0$ (250,1)	7.6 ± 1.0	10.0 ± 1.2
$\tilde{\chi}_1^\pm \tilde{\chi}_1^0$ (500,1)	0.9 ± 0.1	6.3 ± 0.2
$\tilde{\chi}_1^\pm \tilde{\chi}_1^0$ (500,125)	1.0 ± 0.1	5.5 ± 0.2
$\tilde{\chi}_1^\pm \tilde{\chi}_1^0$ (350,100)	2.7 ± 0.3	8.0 ± 0.5

Table 2. Expected and observed event yields in the signal regions. The uncertainties shown include both statistical and systematic sources. The correlation coefficient for the background prediction between the two bins is 0.61. Predicted yields are shown also for several signal models with the masses $(m_{\tilde{\chi}_1^\pm}, m_{\tilde{\chi}_1^0})$ indicated in GeV and with statistical-only uncertainties.

For these models, the kinematic properties of the events are most similar to those from SM backgrounds, and as a result, the acceptance is smaller than for models with larger Δm . For these models with compressed mass spectra, the largest uncertainties in the signal yields arise from the jet energy scale (up to 40%), E_T^{miss} resolution in fast simulation (up to 50%), and limited size of MC samples (up to 60%). These uncertainties reach their maximal values only for models where the acceptance of this analysis is very small and the sensitivity is similarly small. For models with large Δm , where this analysis has the best sensitivity, these uncertainties typically amount to only a few percent. Other experimental and theoretical uncertainties are also considered and lead to small changes in the expected yields. These include effects from the renormalization and factorization scales assumed in the generator on the signal acceptance, the b tagging efficiency, the lepton reconstruction, identification, and isolation efficiency, the trigger efficiency, and the modeling of pileup. Finally, the uncertainty in the integrated luminosity is 2.5% [87].

Figure 6 shows the expected and observed 95% confidence level (CL) exclusion limits for $\tilde{\chi}_1^\pm \tilde{\chi}_2^0 \rightarrow W^\pm H \tilde{\chi}_1^0 \tilde{\chi}_1^0$ as a function of $m_{\tilde{\chi}_1^\pm}$ assuming $m_{\tilde{\chi}_1^0} = 1 \text{ GeV}$ (left) and then in the two-dimensional plane of $m_{\tilde{\chi}_1^\pm}$ and $m_{\tilde{\chi}_1^0}$ (right). This search excludes $m_{\tilde{\chi}_1^\pm}$ values between 220 and 490 GeV when $m_{\tilde{\chi}_1^0} = 1 \text{ GeV}$, and $m_{\tilde{\chi}_1^0}$ values up to 110 GeV when $m_{\tilde{\chi}_1^\pm}$ is around 450 GeV.

Source	Typical range of values [%]
Integrated luminosity	2.5
Size of MC samples	2–60
Pileup	1–5
Renormalization and factorization scales	1–3
ISR modeling	1–5
b tagging efficiency	2–8
Lepton efficiency	2–5
Trigger efficiency	1–5
Jet energy scale	1–40
Fastsim E_T^{miss} resolution	1–50

Table 3. Sources of systematic uncertainty in the estimated signal yield, along with their typical values. The ranges represent variation across the signal masses probed.

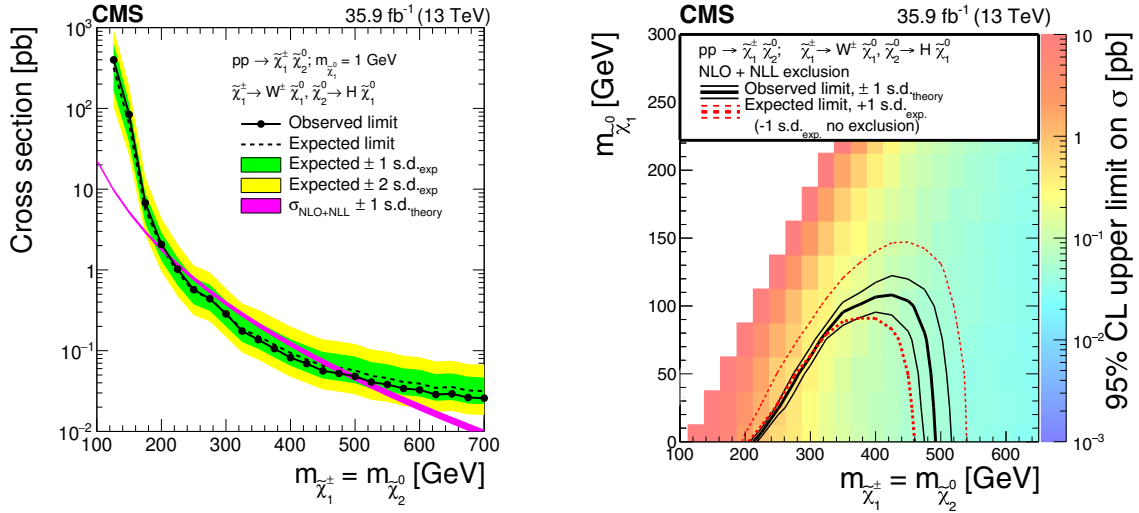


Figure 6. (left) Cross section exclusion limits at the 95% CL are shown for $\tilde{\chi}_1^\pm \tilde{\chi}_2^0 \rightarrow W^\pm H \tilde{\chi}_1^0 \tilde{\chi}_1^0$ as a function of $m_{\tilde{\chi}_1^\pm}$, assuming $m_{\tilde{\chi}_1^0} = 1$ GeV. The solid black line and points represent the observed exclusion. The dashed black line represents the expected exclusion, while the green and yellow bands indicate the ± 1 and 2 standard deviation (s.d.) uncertainties in the expected limit. The magenta line shows the theoretical cross section with its uncertainty. (right) Exclusion limits at the 95% CL in the plane of $m_{\tilde{\chi}_1^\pm}$ and $m_{\tilde{\chi}_2^0}$. The area below the thick black (dashed red) curve represents the observed (expected) exclusion region. The thin dashed red line indicates the $+1 \text{ s.d.}_{\text{exp.}}$ experimental uncertainty. The $-1 \text{ s.d.}_{\text{exp.}}$ line does not appear as no mass points would be excluded in that case. The thin black lines show the effect of the theoretical uncertainties ($\pm 1 \text{ s.d.}_{\text{theory}}$) on the signal cross section.

7 Summary

A search is performed for beyond the standard model physics in events with a leptonically decaying W boson, a Higgs boson decaying to a $b\bar{b}$ pair, and large transverse momentum imbalance. The search uses proton-proton collision data recorded by the CMS experiment in 2016 at $\sqrt{s} = 13$ TeV, corresponding to an integrated luminosity of 35.9 fb^{-1} . The event yields observed in data are consistent with the estimated standard model backgrounds. The results are used to set cross section limits on chargino-neutralino production in a simplified supersymmetric model with degenerate masses for $\tilde{\chi}_1^\pm$ and $\tilde{\chi}_2^0$ and with the decays $\tilde{\chi}_1^\pm \rightarrow W^\pm \tilde{\chi}_1^0$ and $\tilde{\chi}_2^0 \rightarrow H \tilde{\chi}_1^0$. Values of $m_{\tilde{\chi}_1^\pm}$ between 220 and 490 GeV are excluded at 95% confidence level by this search when the $\tilde{\chi}_1^0$ is massless, and values of $m_{\tilde{\chi}_1^0}$ are excluded up to 110 GeV for $m_{\tilde{\chi}_1^\pm} \approx 450$ GeV. These results significantly extend the previous best limits, by up to 270 GeV in $m_{\tilde{\chi}_1^\pm}$ and up to 90 GeV in $m_{\tilde{\chi}_1^0}$.

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